

双曲扁壳式閘門应力計算

陈启強 黃日忠 陈大新 王輝丰 邓健新

§ 1 引言

在一般建筑上应用双曲扁壳是属于簡支而且是均布荷重的情形,但双曲扁壳式閘門则是非簡支非均布荷重的情形,本文任务是对此情形提出应力計算方程及公式,由于我們都是沒有学过薄壳理論,这次是边学边干,所以难免存在很多錯誤,希望各界人士严加指正,我們不胜感謝。

在叙述前先說明本文所采用的符号如下,

A_1, A_2 ——拉梅系数; R_1, R_2 ——曲率半径; ϵ_1, ϵ_2 ——伸长应变; ϵ ——切变; v, v_2 , w ——位移函数; D ——抗变刚度; f_a, f_b ——扁壳的矢高。至于其他符号的意义和一般力学书上的相同。

§ 2 問題假設

一般薄壳問題的方程很复杂,往往对具体問題要作一系列假設,才把方程簡化,当然这些做法要保証在允許的誤差范围之内,現在我們对問題作如下假設:

1. 以扁壳中面的主曲率綫作曲綫坐标系,可近地似用平面笛氏坐标系代替。則有 $ps_1 = A_1 d_1 X$, $ds_2 = A_2 d_2 X$, $\therefore dx \approx d_1 X$, $dy \approx d_2 X$, $\therefore A_1 = A_2 = 1$ 。

2. $S_{12} = T_{12} = T_{21}$ 。

3. $\frac{N_1}{R_1}$, $\frac{N_2}{R_2}$ 可忽略不計。

4. 在曲率改变率 K_1 , K_2 和扭率 L 公式中忽略 $\frac{u}{R_1}$, $\frac{v}{R_2}$, $\frac{1}{R_1} \frac{\partial u}{\partial y}$, $\frac{1}{R_2} \frac{\partial u}{\partial x}$ 。

5. 忽略 $\frac{1}{R_1 R_2}$ 同阶項。

§ 2 建立方程

根据平衡方程、連續性方程、变形方程①以及假設 1—5 得:

$$\frac{\partial T_1}{\partial x} + \frac{\partial s}{\partial y} + q_1 = 0 \quad \dots\dots\dots(1)$$

$$\frac{\partial s}{\partial x} + \frac{\partial T_2}{\partial y} + q_2 = 0 \quad \dots\dots\dots(2)$$

$$\frac{\partial M_1}{\partial x} + \frac{\partial M_{21}}{\partial y} - N_1 = 0 \quad \dots\dots\dots(3)$$

$$\frac{\partial M_1}{\partial x} + \frac{\partial M_2}{\partial y} - N_2 = 0 \quad \dots\dots\dots(4)$$

$$\frac{K_1}{R_2} + \frac{K_2}{R_1} + \frac{\partial^2 \varepsilon_2}{\partial x^2} + \frac{\partial^2 \varepsilon_1}{\partial y^2} - \frac{\partial^2 \omega}{\partial x \partial y} = 0 \quad \dots\dots\dots(5)$$

$$K_1 = -\frac{\partial^2 w}{\partial x^2}, \quad K_2 = -\frac{\partial^2 w}{\partial y^2}, \quad L = -\frac{\partial^2 w}{\partial x \partial y}$$

另外从薄壳理論还知道: $M_1 = D(k_1 + \mu k_2)$, $M_{21} = \frac{E\delta_3}{12(1+\mu)} L$.

則 $M_1 = -D\left(\frac{\partial^2 w}{\partial x^2} + \mu \frac{\partial^2 w}{\partial y^2}\right)$, $M_{21} = -D(1-\mu)\frac{\partial^2 w}{\partial x \partial y}$, M_1, M_{21} 代入(3)(4)得:

$N_1 = -D\frac{\partial}{\partial x}(\Delta w)$, $N_2 = -D\frac{\partial}{\partial y}(\Delta w)$, 又将 N_1, N_2 代入平衡方程之三式①得:

$$\frac{T_1}{R_1} + \frac{T_2}{R_2} + D\Delta\Delta(w) = q_n \quad \dots\dots\dots(6)$$

从問題的提法知 $q_1 = 0$, 而 q_2 是自重所产生的力, 而 $q_2^* = -r$ (单位体重) q_n 则是水压力, 設水面至薄壳上頂之高度为 h , 則 $q_n = \gamma + h$.

引入应力函数 φ 并使 $T_1 = \frac{\partial^2 \varphi}{\partial y^2}$, $T_2 = \frac{\partial^2 \varphi}{\partial x^2} - \gamma_y$, $S = T_{21} = T_{12} = -\frac{\partial^2 \varphi}{\partial x \partial y}$. 显然

滿足(1)(2), 而(6)变为: $\Delta_k \varphi + D\Delta\Delta W = h + \gamma + \frac{\gamma_y}{R_2} \dots\dots\dots(7)$

从而中面的变形和内力的关系如下:

$$\left. \begin{aligned} \varepsilon_1 &= \frac{1}{E\delta} (T_1 - \mu T_2) = \frac{1}{E\delta} \left(\frac{\partial^2 \varphi}{\partial y^2} - \mu \frac{\partial^2 \varphi}{\partial x^2} + \mu \gamma_y \right) \\ \varepsilon_2 &= \frac{1}{E\delta} (T_2 - \mu T_1) = \frac{1}{E\delta} \left(\frac{\partial^2 \varphi}{\partial x^2} - \mu \frac{\partial^2 \varphi}{\partial y^2} - \gamma_y \right) \end{aligned} \right\} \dots\dots\dots(8)$$

$$\omega = \frac{2(1+\mu)}{F\delta} S = -\frac{2(1+\mu)}{E\delta} \frac{\partial^2 \varphi}{\partial x \partial y}$$

从而有 $\frac{\partial^2 \varepsilon_1}{\partial y^2} + \frac{\partial^2 \varepsilon_2}{\partial x^2} - \frac{\partial^2 \omega}{\partial x \partial y} = \frac{1}{E\delta} \Delta\Delta\varphi \dots\dots\dots(9)$, 將(8)(9)代入(5)得:

$$\frac{1}{E\delta} \Delta\Delta\varphi - \Delta_k w = 0 \dots\dots\dots(10)$$

現得所得之方程列于下:

$$\begin{cases} \Delta_k \varphi + D\Delta\Delta w = h + \gamma + \frac{\gamma_y}{R_2} \quad \dots\dots\dots(7) \\ \Delta\Delta\varphi - E\delta\Delta_k w = 0 \quad \dots\dots\dots(10) \end{cases}$$

則应力由下列公式求出:

$$T_1 = \frac{\partial^2 \varphi}{\partial y^2}, \quad T_2 = \frac{\partial^2 \varphi}{\partial x^2} - \gamma_y, \quad J = T_{21} = T_{12} = -\frac{\partial^2 \varphi}{\partial x \partial y}, \quad M_1 = -D\left(\frac{\partial^2 w}{\partial x^2} + \mu \frac{\partial^2 w}{\partial y^2}\right),$$

$$M_2 = -D\left(\frac{\partial^2 w}{\partial y^2} + \mu \frac{\partial^2 w}{\partial x^2}\right), \quad M = -\frac{E\delta^3}{12(1+\mu)} \frac{\partial^2 w}{\partial x \partial y}$$

上面 $\Delta_k = \frac{1}{R_1} \frac{\partial^2}{\partial y^2} + \frac{1}{R_2} \frac{\partial^2}{\partial x^2}$; $\Delta = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2}$; γ ——单位体积壳体重量。

§ 边界条件

1. 支承情况 双曲扁壳是垂直放在水底, 故一边受重力作用。而曲面受到垂直于面的水压力, 为了简单起见不计算水对四边的压力, 此外, 还有一个问题要考虑, 就是 $x=0$, $x=a$ 时, 由于闸门的自重作用在 oy 方向产生位移 v , 为了使 $v=0$, 故必须在边界上有一应力 $T_2 = -Yy$, 抵消自重的作用。

2. 边界条件

a) $x=0$, $x=a$, $T_1 = M_1 = V = W = 0$, $T_2 = -Yy$, 则可推得 $\frac{\partial^2 \varphi}{\partial y^2} = 0$,

$$\frac{\partial^2 \varphi}{\partial x^2} = 0, \quad \frac{\partial^2 w}{\partial x^2} = -\mu \frac{\partial^2 w}{\partial y^2}.$$

b) $y=0$, $T_2 = M_2 = U = W = 0$, 这样可得 $\frac{\partial^2 \varphi}{\partial x^2} = 0$, $\frac{\partial^2 \varphi}{\partial y^2} = 0$,

$$\frac{\partial^2 w}{\partial y^2} = -\mu \frac{\partial^2 w}{\partial x^2}.$$

c) $y=b$, $M_2 = u = w = 0$, $T_2 = -Yb$, 同样得 $\frac{\partial^2 \varphi}{\partial x^2} = 0$, $\frac{\partial^2 \varphi}{\partial y^2} = 0$,

$$\frac{\partial^2 w}{\partial y^2} = -\mu \frac{\partial^2 w}{\partial x^2}.$$

§ 3 问题解法

用双三角级数求解, 令:

$$\varphi(x, y) = \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} B_{mn} \sin \frac{m\pi}{a} x \sin \frac{n\pi}{b} y \dots\dots\dots (2)$$

$$w(x, y) = \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} A_{mn} \sin \frac{m\pi}{a} x \sin \frac{n\pi}{b} y \dots\dots\dots (3)$$

显然满足所有的边界条件。

$$\text{设 } L_{mn}^{(1)} = \left[\left(\frac{m\pi}{a} \right)^2 + \left(\frac{n\pi}{b} \right)^2 \right]^2, \quad L_{mn}^{(2)} = \left[\frac{1}{R_1} \left(\frac{n\pi}{b} \right)^2 + \frac{1}{R_2} \left(\frac{m\pi}{a} \right)^2 \right], \text{ 将 (2)}$$

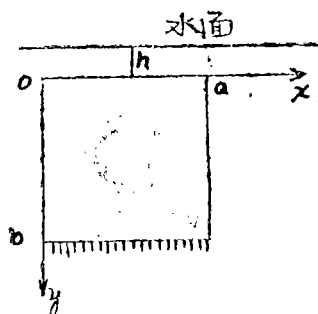
$$(3) \text{ 代入 (1) 之第二式得 } B_{mn} = \frac{-E\delta L_{mn}^{(2)} A_{mn}}{L_{mn}^{(1)}} \dots\dots\dots (4)$$

又将 (1) 之第一式的右边展成重三角级数:

$$h + y + \frac{\gamma y}{\rho_2} = \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} c_{mn} \sin \frac{m\pi}{a} x \sin \frac{n\pi}{b} y \dots\dots\dots (5)$$

$$\text{可求出 } c_{2m+2n} = c_{2m+2n-1} = 0; \quad c_{2m-1, 2n} = \frac{-4b(R_2 + \gamma)}{R_2 n (2m-1) \pi^2}; \quad c_{2m-1, 2n-1} = \frac{16R_2 b + 8b(R_2 + \gamma)}{(2m-1)(2n-1)R_2 \pi^2}$$

$$\text{将 (2)(3)(5) 代入 (1) 一式得: } DL_{mn}^{(1)} A_{mn} - L_{mn}^{(1)} B_{mn} = c_{mn} \dots\dots\dots (6)$$



$$\therefore \text{从(1)(6)得: } A_{mn} = \frac{L_{mn}^{(1)} c_{1mn}}{DL_{mn}^{(1)^2} + E\delta L_{mn}^{(2)^2}} \dots\dots\dots (7)$$

$$B_{mn} = \frac{-E\delta L_{mn}^{(2)} c_{1mn}}{DL_{mn}^{(1)^2} + E\delta L_{mn}^{(2)^2}} \dots\dots\dots (8)$$

这样应力依下列公式求出:

$$T_1 = \frac{\partial^2 \varphi}{\partial y^2} = -\left(\frac{\pi}{b}\right)^2 \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} n^2 B_{mn} \sin \frac{m\pi}{a} x \sin \frac{n\pi}{b} y \dots\dots\dots (9)$$

$$T_2 = \frac{\partial^2 \varphi}{\partial x^2} - \gamma y = -\left(\frac{\pi}{a}\right)^2 \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} m^2 B_{mn} \sin \frac{m\pi}{a} x \sin \frac{n\pi}{b} y - \gamma y \dots\dots\dots (10)$$

$$S = -\frac{\partial^2 \varphi}{\partial x \partial y} = -\frac{\pi^2}{ab} \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} mn A_{mn} \cos \frac{m\pi}{a} x \cos \frac{n\pi}{b} y \dots\dots\dots (11)$$

$$M_1 = -D \left(\frac{\partial^2 w}{\partial x^2} + \mu \frac{\partial^2 w}{\partial y^2} \right) = \frac{D\pi^2}{a^2 b^2} \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} A_{mn} (b^2 m^2 + \mu a^2 n^2) \sin \frac{m\pi}{a} x \sin \frac{n\pi}{b} y \dots\dots (12)$$

$$M_2 = -D \left(\frac{\partial^2 w}{\partial y^2} + \mu \frac{\partial^2 w}{\partial x^2} \right) = \frac{D\pi^2}{a^2 b^2} \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} A_{mn} (a^2 n^2 + \mu b^2 m^2) \sin \frac{m\pi}{a} x \sin \frac{n\pi}{b} y \dots\dots (13)$$

$$U = \frac{E\delta^3}{12(1+\mu)} \frac{\partial^2 w}{\partial x \partial y} = -\frac{D\pi}{ab} \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} mn A_{mn} \cos \frac{m\pi}{a} x \cos \frac{n\pi}{b} y \dots\dots\dots (14)$$

公式中 A_{mn} 及 B_{mn} 由 (7) (8) 决定。

双曲扁壳式闸门的具體数据如下:

$$a = 2.9\text{m}, \quad b = 3.4\text{m}, \quad fa = fb = 0.5\text{m}, \quad \delta = 10\text{cm}, \quad E = 38 \times 10^4 \text{kg/cm}^2, \quad \mu = 0.2$$

$$\text{曲面方程: } z = fa + fb - \frac{4fa}{a^2} \left(\frac{a}{2} - x \right)^2 - \frac{4fb}{b^2} \left(\frac{b}{2} - y \right)^2; \quad k_1 = \frac{1}{R_1} = \frac{3fa}{a^2}, \quad k_2 = \frac{1}{R_2} = \frac{3fb}{b^2}$$

$$\gamma = \frac{E\delta^3}{12(1+\mu^2)} \text{ 将这些数据代入 (9) } \dots\dots (14) \text{ 并取 } m = 1, 3; \quad n = 1, 2 \text{ 得:}$$

$$T_1 = 0.853 \left(8201.231 \sin \frac{\pi}{2.9} x \sin \frac{\pi}{3.4} y - 2098.636 \sin \frac{\pi}{2.9} x \sin \frac{2\pi}{3.4} y \right. \\ \left. + 293.271 \sin \frac{3\pi}{2.9} x \sin \frac{\pi}{3.4} y - 31.704 \sin \frac{3\pi}{2.9} x \sin \frac{2\pi}{3.4} y \right)$$

$$T_2 = 1.173 \left(8201.231 \sin \frac{\pi}{2.9} x \sin \frac{\pi}{3.4} y - 524.659 \sin \frac{\pi}{2.9} x \sin \frac{2\pi}{3.4} y \right. \\ \left. + 2694.139 \sin \frac{3\pi}{2.9} x \sin \frac{\pi}{3.4} y - 133.834 \sin \frac{3\pi}{2.9} x \sin \frac{2\pi}{3.4} y \right) - 2400y$$

$$S = 8201.231 \cos \frac{\pi}{2.9} x \cos \frac{\pi}{3.4} y - 1049.317 \cos \frac{\pi}{2.9} x \cos \frac{2\pi}{3.4} y \\ + 394.814 \cos \frac{3\pi}{2.9} x \cos \frac{\pi}{3.4} y - 122.556 \cos \frac{3\pi}{2.9} x \cos \frac{2\pi}{3.4} y$$

$$\begin{aligned}
 M_1 &= 44.235 \times 10^{-2} \left(12706.37 \sin \frac{\pi}{2.9} x \cos \frac{\pi}{3.4} y - 1671.507 \sin \frac{\pi}{2.9} x \sin \frac{2\pi}{3.4} y \right. \\
 &\quad \left. + 22195.185 \sin \frac{3\pi}{2.9} x \sin \frac{\pi}{3.4} y - 2076.534 \sin \frac{3\pi}{2.9} x \sin \frac{2\pi}{3.4} y \right) \\
 M_2 &= 44.235 \times 10^{-2} \left(9243.994 \sin \frac{\pi}{2.9} x \sin \frac{\pi}{3.4} y - 4864.142 \sin \frac{\pi}{2.9} x \sin \frac{2\pi}{3.4} y \right. \\
 &\quad \left. + 1794.13 \sin \frac{3\pi}{2.9} x \sin \frac{\pi}{3.4} y - 671.421 \sin \frac{3\pi}{2.9} x \sin \frac{2\pi}{3.4} y \right) \\
 H &= 3.17 \times 10^{-2} \left(-1099.167 \cos \frac{\pi}{2.9} x \cos \frac{\pi}{3.4} y + 289.118 \cos \frac{\pi}{2.9} x \cos \frac{\pi}{3.4} y \right. \\
 &\quad \left. - 639.999 \cos \frac{3\pi}{2.9} x \cos \frac{\pi}{3.4} y + 119.754 \cos \frac{3\pi}{2.9} x \cos \frac{2\pi}{3.4} y \right)
 \end{aligned}$$

§ 5 应力分布情况

点 应力	(0,0)	(0, $\frac{b}{2}$)	(0,b)	($\frac{a}{2}$,0)	(a,0)	(a, $\frac{b}{2}$)	($\frac{a}{2}$,b)	(a,b)
T_1	0	0	0	0	0	0	0	0
T_2	0	-4.408	-8.160	0	0	-4.080	-8.160	-8.160
S	7.923	1.192	-10.263	0	-7.923	-1.192	0	10.268

($\frac{3}{4}a, \frac{b}{2}$)	($\frac{a}{2}, \frac{3}{4}b$)	($\frac{a}{4}, \frac{b}{4}$)	($\frac{a}{2}, \frac{b}{2}$)	($\frac{a}{4}, \frac{b}{2}$)	($\frac{a}{2}, \frac{b}{4}$)	($\frac{3}{4}a, \frac{3}{4}b$)
5.100	6.484	2.302	6.725	5.100	4.975	4.940
4.955	-3.121	3.757	2.374	4.955	2.091	0.852
-0.655	0	3.653	0	0.655	0	3.653

表中应力单位: T/m²

①见 薄壳理论

Calculation for the stress of double curved flat-shell type water-gate

Faculty of Mathematic-Mechanic Class four

Huang ri-zhong Chen da-xin

Wang hui-feng deng jian-xin

There are broad application of hyperbolic flat-shell in civil engineering but most of them are of simple condition of simple support and distributed load. Now the double curved flat-shell type water-gate is neither simple support nor distributed load. This article propose the equations and for mule for calculating this problem in cluding the following contents

1. hypothesis problem
2. equation established
3. boundary condition
4. Calcululng formule and resuet